

Heat capacity and thermal expansion coefficient of rare earth uranates $\text{RE}_6\text{UO}_{12}$ (RE = Nd, Gd and Eu)

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Abstract Rare earth uranates $\text{Nd}_6\text{UO}_{12}$, $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$ were prepared by combustion synthesis and characterized by XRD. Single-phase rhombohedral structure was observed for all the above compounds. Heat capacity measurements were carried out on $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ with differential scanning calorimetry in the temperature range 298–800 K. Enthalpy, entropy and Gibbs energy functions were computed. Heat capacity values of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ at 298 K are 436 and 400 $\text{J K}^{-1} \text{mol}^{-1}$, respectively. Thermal expansion characteristics were studied using high temperature X-ray diffraction (HTXRD) in the temperature range 298–873 K. The coefficients of thermal expansion measured for $\text{Eu}_6\text{UO}_{12}$ are 10.5×10^{-6} and $7.3 \times 10^{-6} \text{K}^{-1}$ along *a*- and *c*-axis, respectively. Similarly, the coefficients of thermal expansion of $\text{Gd}_6\text{UO}_{12}$ along *a*-axis are $10.0 \times 10^{-6} \text{K}^{-1}$ and along *c*-axis is $9.7 \times 10^{-6} \text{K}^{-1}$.

Keywords Heat capacity · Rare earths · Uranates · Rare earth uranates · Thermal expansion

Introduction

$\text{RE}_6\text{UO}_{12}$ (RE = rare earth) type of compounds are known to exist in RE–U–O system [1, 2] with rhombohedral structure [3–5]. Though $\text{RE}_6\text{UO}_{12}$ compounds are not formed under normal reactor operation conditions, they might be formed

under transient conditions and therefore thermophysical properties such as heat capacity and thermal expansion of these compounds are of interest. Thermal expansion of $\text{RE}_6\text{UO}_{12}$ (RE = La, Nd and Sm,) were measured by Jena et al. [5, 6] and thermal diffusivity of $\text{RE}_6\text{UO}_{12}$ (RE = La, Gd, Dy) by Krishnaiah et al. [7]. Differential scanning calorimetry is widely used to measure heat capacity of variety of compounds such as ceramics, organic materials, polymers, ionic liquids, etc., [8–17]. Heat capacity of $\text{RE}_6\text{UO}_{12}$ (RE = La, Sm and Eu) had been measured by DSC in our previous study [18]. In the present study heat capacity of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ were measured by DSC and thermal expansion of $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$ by HTXRD.

Experimental

Sample preparation

UO_2 of nuclear grade purity supplied by Nuclear Fuel Complex, Hyderabad, India, and RE_2O_3 (RE = Nd, Gd and Eu) of 99.9% purity supplied by M/s. Indian Rare Earths were used for preparing the samples. $\text{RE}_6\text{UO}_{12}$ compounds were prepared by combustion synthesis, as explained in our previous paper [18]. Stoichiometric ratio of oxidant to fuel was used for the preparation. The sample pellets obtained were heated at 1,473 K in air for 8 h. As the samples were heated in air, the $\text{RE}_6\text{UO}_{12}$ compounds prepared are expected to be stoichiometric [5]. Since citric acid was used as the fuel, the samples were analyzed for carbon impurity by using carbon analyzer of M/s. Eltra, Germany (ELTRACS 800). All the samples were found to contain <100 ppm of carbon. The sample pellets were then ground and X-ray diffraction analysis was performed on all the above mentioned samples, using Siemens D500 powder

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diffractometer using $\text{CuK}\alpha$ radiation monochromatized with curved graphite crystal monochromator placed in front of the NaI(Tl) scintillation detector. The scans were made in the range of $10^\circ \leq 2\theta \leq 80^\circ$. The X-ray diffraction patterns of $\text{Nd}_6\text{UO}_{12}$, $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$ are shown in Fig. 1, which are in good agreement with the reported literature patterns (JCPDS Card No.190830, JCPDS Card No. 801372, JCPDS Card No. 201335). The computed lattice parameter values of $\text{RE}_6\text{UO}_{12}$ listed in Table 1 are in good agreement with the values reported in the literature [3, 6] (JCPDS Card No. 201335). The lattice parameter values also suggest that the $\text{RE}_6\text{UO}_{12}$ compounds formed are stoichiometric. For DSC measurements, the powders of these samples were compacted into pellets of 5 mm in diameter and sintered at 1,673 K in air for 8 h.

Heat capacity measurements

A heat flux type differential scanning calorimeter (model number DSC821e/700 of M/s. Mettler Toledo GmbH,

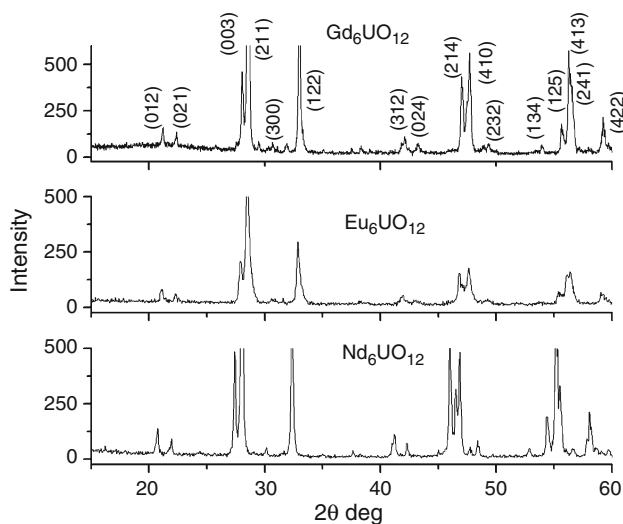


Fig. 1 Room temperature XRD patterns of $\text{RE}_6\text{UO}_{12}$

Switzerland) was used in this study. Pellet samples were hermetically sealed in 40 μL Al-pans and a pinhole was made on the lid. High purity argon was used as the purge gas at a flow rate of 50 mL min^{-1} and a heating rate of 10 K min^{-1} was employed in DSC measurements in the temperature range 298–800 K. A thin disc of sapphire was used as the heat capacity standard. A three segment-heating programme was used. The first segment lasting for 5 min was an isothermal one at the initial temperature; the second segment was a dynamic one with a heating rate of 10 K min^{-1} and the final segment lasting for 5 min was another isothermal one at the final temperature. The “Classical” three step procedure of three runs was used for measuring heat capacity. The first run was a blank run with empty pans on both the sample and the reference sides of the calorimeter. The blank run was for the correction of any asymmetry between the two pans. The second run was the calibration run where the calibration material, a disc of sapphire, was placed on the sample pan whereas the pan on the reference side was kept empty. The third run was the sample run where the sample in the form of pellet was placed on the sample pan with empty pan on the reference side. These three runs were repeated for each measurement of heat capacity. Temperature, heat and heat rate calibrations of the DSC were carried out, as explained in our earlier study [19, 20].

Thermal expansion measurements by HTXRD

HTXRD measurements were carried out by using the HTK-16 module of Philips X’Pert Pro MPD system. The powder samples were mixed with Pt powder (1:5 = Pt: $\text{RE}_6\text{UO}_{12}$ wt. ratio) as internal standard for temperature correction. The fine powders were mounted on the tantalum strip (sample carrier) that was resistively heated at programmed heating rates. The sample temperature was monitored with a thermocouple spot-welded to the rear side of the sample carrier strip. All measurements were carried out in a

Table 1 Crystal data for the $\text{RE}_6\text{UO}_{12}$ compounds measured by XRD

Crystal data at 298 K		Compound					
		Space group- $R\bar{3}$					
		$\text{Nd}_6\text{UO}_{12}$		$\text{Gd}_6\text{UO}_{12}$		$\text{Eu}_6\text{UO}_{12}$	
		This study	Literature values	This study	Literature values	This study	Literature values
Hexagonal parameters	a (\AA)	10.254 (1)	10.254 [12]	10.098 (2)	10.070 [11] 10.077 [17]	10.140 (4)	10.126 [6, 12]
	c (\AA)	9.748 (1)	9.748 [12]	9.554 (2)	9.529 [11] 9.526 [17]	9.569 (3)	9.601 [6, 12]
Rhombohedral parameters	a_R (\AA)	6.754 (3)	–	6.643 (2)	6.6285 [17]	6.649 (4)	6.665 [6, 12]
	α (deg)	98.79	–	98.89	98.94 [17]	98.95	98.87 [6, 12]

vacuum of about 10^{-5} Torr. The XRD patterns were recorded in step scan mode with the step size of $0.02^\circ \text{ s}^{-1}$. The XRD patterns of the samples were recorded at room temperatures, and then the temperature was increased to 873 K in steps of 100 K. The lattice parameters were obtained from the XRD data at each temperature. The lattice parameters, l , were fitted against temperature (in Kelvin) to the polynomial expression: $l = x + yT + zT^2$. The average axial expansion coefficient, α_l , between the temperatures, T_1 and T_2 , was calculated by the equation, $\alpha_l = (l_2 - l_1)/(l_1(T_2 - T_1))$, where l_1 and l_2 are the lattice parameters at temperatures T_1 and T_2 , respectively. The cell volume for this hexagonal system was computed by the expression $v = \sqrt{3}/2(a^2c)$ and the percentage variation in unit cell volume was computed by the expression $((v_T - v_{298})/v_{298}) \times 100$.

Results

Heat capacity of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$

Heat capacity data of sapphire given by National Institute of Standards and Technology, USA (NIST) were used for computing the heat capacities of the samples. Heat capacity of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ measured by DSC and given in Tables 2 and 3, respectively, are the mean of eight

measurements. The error in the heat capacity values by our DSC measurements had been earlier determined to be $\pm 3\%$ by measurements on ThO_2 samples [21]. However, the relative standard deviations among the present measurements are in the range of 1–3%. Hence the uncertainty in the measured heat capacity values has been estimated to be $\pm 3\%$. The measured heat capacity of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ were least squares fitted to obtain the following polynomial in temperature.

$$C_{p,m}(\text{Nd}_6\text{UO}_{12})(\text{JK}^{-1} \text{ mol}^{-1}) = 702.196 - 1.3175(T/K) - 5.011332 \times 10^6 (K/T)^2 + 2.42 \times 10^{-3}(T/K)^2 - 1.3314 \times 10^{-6} (T/K)^3 \times (298 - 800\text{K}) \quad (1)$$

$$C_{p,m}(\text{Gd}_6\text{UO}_{12})(\text{JK}^{-1} \text{ mol}^{-1}) = 500.68 + 7.07 \times 10^{-3} (T/K) - 9.161807 \times 10^6 (K/T)^2 (298 - 800\text{K}) \quad (2)$$

The standard errors of the fit equations 1 and 2 are 2.8 and $2.5 \text{ J K}^{-1} \text{ mol}^{-1}$, respectively. The measured heat capacity of $\text{Nd}_6\text{UO}_{12}$, and $\text{Gd}_6\text{UO}_{12}$ along with the fit values are shown in Figs. 2 and 3, respectively. The heat capacity of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ computed by Neumann–Kopp’s law using the heat capacity data of UO_3 [22], Nd_2O_3 [23] and Gd_2O_3 [23] are also shown in Figs. 2 and 3, respectively. The heat capacity data of

Table 2 Thermodynamic functions of $\text{Nd}_6\text{UO}_{12}$

T/K	$C_{p,m}/\text{J K}^{-1} \text{ mol}^{-1}$		$H_T^0 - H_{298}^0/\text{J mol}^{-1}$	$S_T^0/\text{J K}^{-1} \text{ mol}^{-1}$	$G_T^0 - H_{298}^0/T/\text{J K}^{-1} \text{ mol}^{-1}$
	Measured	Fit			
298	436	433	0	574	−574
300	436	433	866	578	−575
400	448	446	44,820	729	−617
500	458	462	90,173	841	−661
600	489	481	137,335	933	−704
700	497	499	186,395	1,013	−746
800	512	508	236,821	1,082	−786

Table 3 Thermodynamic functions of $\text{Gd}_6\text{UO}_{12}$

T/K	$C_{p,m}/\text{J K}^{-1} \text{ mol}^{-1}$			$H_T^0 - H_{298}^0/\text{J mol}^{-1}$	$S_T^0/\text{J K}^{-1} \text{ mol}^{-1}$	$G_T^0 - H_{298}^0/T/\text{J K}^{-1} \text{ mol}^{-1}$
	Measured	Fit	Sahu et al. [17]			
298	400	400	386	0	551	−551
300	401	401	386	801	553	−551
400	446	446	416	43,481	676	−567
500	471	468	439	89,286	778	−599
600	480	479	454	136,689	864	−637
700	483	487	459	185,036	939	−675
800	499	492	456	233,998	1,004	−712

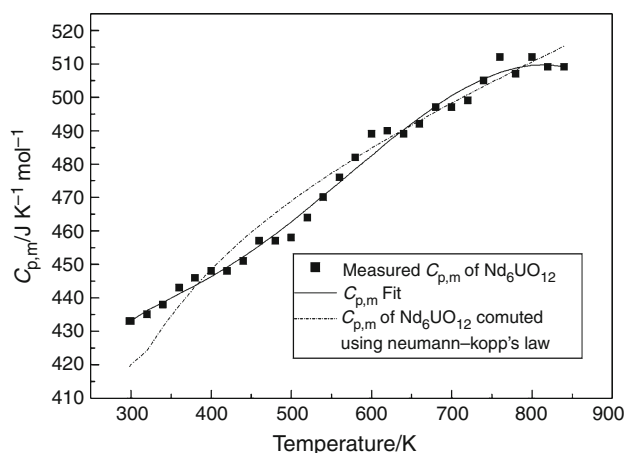


Fig. 2 Heat capacity data of $\text{Nd}_6\text{UO}_{12}$

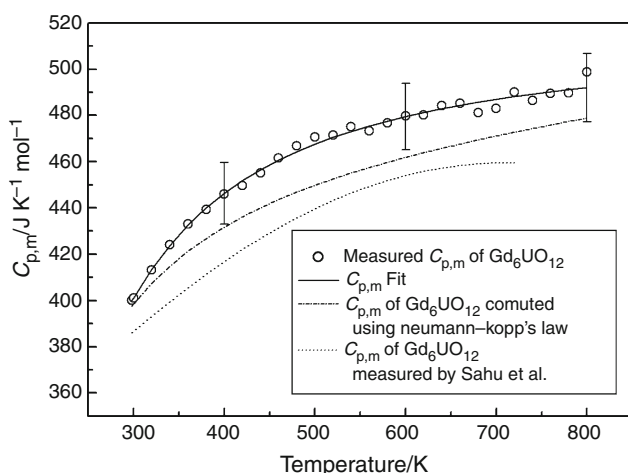


Fig. 3 Heat capacity data of $\text{Gd}_6\text{UO}_{12}$

respective component oxides used for the computation are given below

$$C_{p,m}(\text{UO}_3)(\text{J K}^{-1} \text{mol}^{-1}) = 90.374 + 11.046 \times 10^{-3} (T/\text{K}) - 1.109 \times 10^6 (K/T)^2 \quad (3)$$

$$C_{p,m}(\text{Nd}_2\text{O}_3)(\text{J K}^{-1} \text{mol}^{-1}) = 115.52 + 29.999 \times 10^{-3} (T/\text{K}) - 1.167 \times 10^6 (K/T)^2 \quad (4)$$

$$C_{p,m}(\text{Gd}_2\text{O}_3)(\text{J K}^{-1} \text{mol}^{-1}) = 119.206 + 12.945 \times 10^{-3} (T/\text{K}) - 1.561 \times 10^6 (K/T)^2 \quad (5)$$

As can be seen in the figure the present data are in good agreement within $\pm 3\%$ with that computed using Neumann–Kopp's law using the heat capacity of their respective component oxides. However, the heat capacity data of $\text{Gd}_6\text{UO}_{12}$ measured by Sahu et al. [24] are about 3–6% less than the present data. From the heat capacity data, other thermodynamic functions such as enthalpy, entropy and Gibbs energy functions of $\text{Nd}_6\text{UO}_{12}$ and

$\text{Gd}_6\text{UO}_{12}$ were computed and are given in Tables 2 and 3, respectively. The S_{298}^0 values of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ required for the computation of entropies were estimated by using S_{298}^0 of their respective component oxides (Nd_2O_3 , Gd_2O_3 and UO_3) from the literature [22, 23] as explained below. The estimated value of S_{298}^0 for $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ are 574.5 and 550.7 $\text{J K}^{-1} \text{mol}^{-1}$, respectively.

$$S_{298}^0(\text{RE}_6\text{UO}_{12}) = 3 S_{298}^0(\text{RE}_2\text{O}_3) + S_{298}^0(\text{UO}_3) \quad (\text{RE} = \text{Gd}, \text{Nd})$$

Thermal expansion of $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$

The lattice parameters of $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$ were estimated using NBSAIDS83 software. The diffraction pattern of the compounds could be indexed on $R\bar{3}$ space group as reported by Hinatsu et al. [2] and Jena et al. [5, 6]. The structural refinement of the samples was carried out by using the structure of iso-structural compounds $\text{La}_6\text{UO}_{12}$ reported by Hinatsu et al. [2]. All reflections were indexed based on hexagonal crystallographic axes; for the Miller indices, the relationship $-h + k + l = 3n$ was observable. The lattice constants were found to expand along a and c axis. The thermal expansion measurements of the samples are presented up to 873 K to compare the effect of it on the heat capacity values of the compounds in the same temperature range. The percentage variation along the c -axis is observed to be comparatively lower than the a -axis. Similar observation was reported in our previous work [5] for $\text{RE}_6\text{UO}_{12}$ ($\text{RE} = \text{La}, \text{Nd}$ and Sm). The observed expansion behaviour can be explained based on crystal structure and bonding of the compounds. The chemical bonding in these compounds can be considered to be ionic in nature. In $\text{Gd}_6\text{UO}_{12}$, six oxygen atoms surround a uranium ion forming a distorted octahedron. Similarly, six oxygen atoms surround the Gd^{3+} ion in a highly distorted octahedral fashion. Hence both Gd and U are in six coordination sphere in $\text{Gd}_6\text{UO}_{12}$. The corners of UO_6 octahedra are shared by GdO_6 octahedra. Two edge shared GdO_6 octahedra separate two UO_6 along c -axis; this gives rigidity along the c -axis and therefore lowers percentage expansion with temperature. The lattice constants vs. T (K) of $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$ are given in Fig. 4. The thermal expansions of the compounds showed an increasing trend with increasing temperature. The lattice constants of $\text{Eu}_6\text{UO}_{12}$ are higher than $\text{Gd}_6\text{UO}_{12}$ exhibiting the trend of lanthanide contraction. The coefficient of thermal expansion measured for $\text{Eu}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ are shown in Table 4. The trends in the percentage variation along a and c axis for $\text{RE}_6\text{UO}_{12}$ computed from the present investigation as well as from our previous study [5, 6] are given in Figs. 5 and 6, respectively. The trends in the variation of

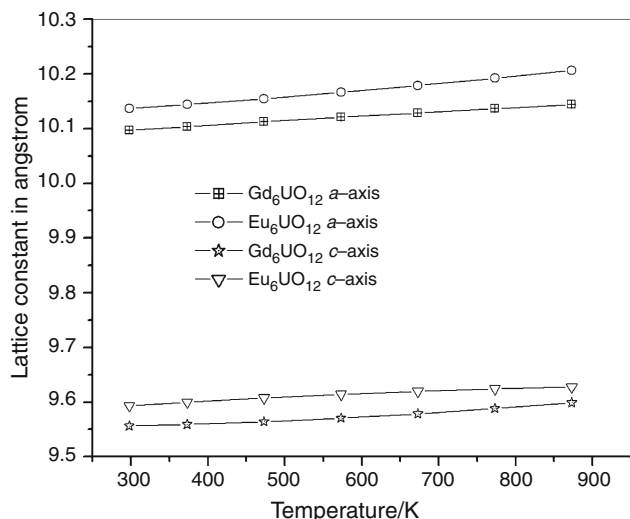


Fig. 4 Lattice expansion along *a*- and *c*-axis of $\text{Gd}_6\text{UO}_{12}$ and $\text{Eu}_6\text{UO}_{12}$

Table 4 The average linear expansion coefficient along the hexagonal axes in the temperature range 298–873 K for $\text{RE}_6\text{UO}_{12}$ compounds

Compounds	$\alpha_a/10^6 \text{ K}^{-1}$	$\alpha_c/10^6 \text{ K}^{-1}$
$\text{Eu}_6\text{UO}_{12}$	10.5 ± 1.0	7.3 ± 1.0
$\text{Gd}_6\text{UO}_{12}$	10.0 ± 1.0	9.7 ± 1.0
$\text{Gd}_6\text{UO}_{12}$ [17]	9.6 ± 1.2	12.3 ± 3.6

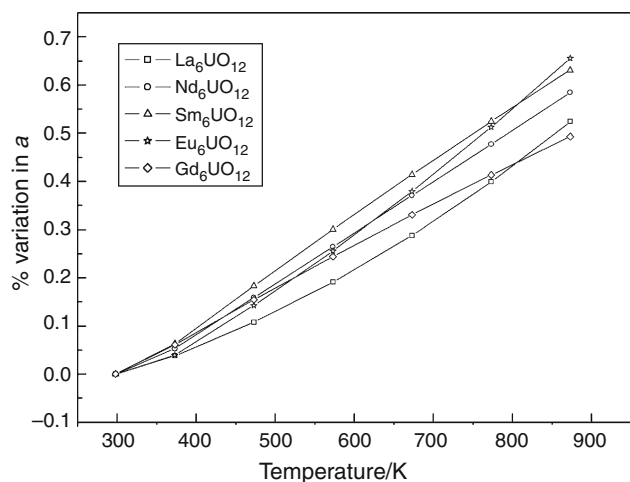


Fig. 5 Trend in % variation along *a*-axis with temperature for $\text{RE}_6\text{UO}_{12}$

percentage volume thermal expansion computed for $\text{RE}_6\text{UO}_{12}$ from the present investigation as well as from our previous study [5, 6] are shown in Fig. 7. As can be seen in the Fig. 7, the percentage volume thermal expansion for $\text{RE}_6\text{UO}_{12}$ varies as $\text{RE} = \text{Sm} > \text{Nd} > \text{Eu} > \text{La} \approx \text{Gd}$.

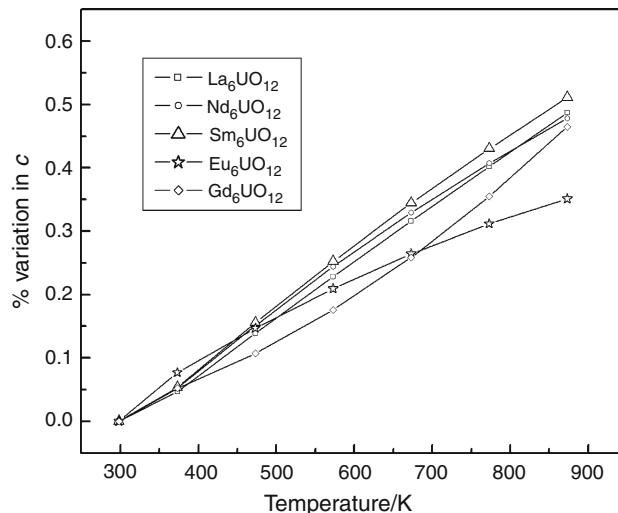


Fig. 6 Trend in % variation along *c*-axis with temperature for $\text{RE}_6\text{UO}_{12}$

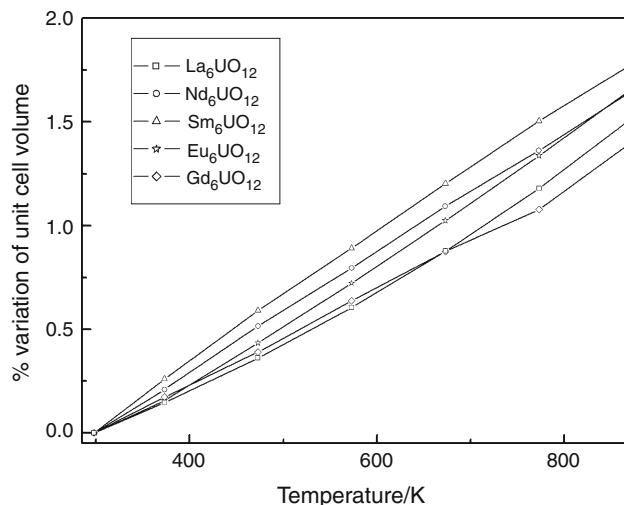


Fig. 7 Variation of percentage volume expansion of $\text{RE}_6\text{UO}_{12}$ with temperature

Discussion

Heat capacity data of $\text{RE}_6\text{UO}_{12}$ ($\text{RE} = \text{Nd}, \text{Gd}$) are compared with that for $\text{RE} = \text{La}, \text{Sm}$ and Eu measured in our previous study [8] in Fig. 8. As can be seen in the Fig. 8, the heat capacity of $\text{RE}_6\text{UO}_{12}$ varies in the order $\text{Eu}_6\text{UO}_{12}$ [8] > $\text{Sm}_6\text{UO}_{12}$ [8] > $\text{Nd}_6\text{UO}_{12}$ > $\text{Gd}_6\text{UO}_{12}$ > $\text{La}_6\text{UO}_{12}$ [18]. Heat capacity of rare earth compounds includes contribution due to (1) phonons arising due to harmonic lattice vibration, (2) dilatational term arising due to thermal expansion and (3) Schottky contribution arising due to promotion of electrons to higher energy levels [24, 25]. As per theoretical expectations the contribution of phonons to heat capacity vary in $\text{RE}_6\text{UO}_{12}$ as $\text{La}^{3+} > \text{Nd}^{3+} >$

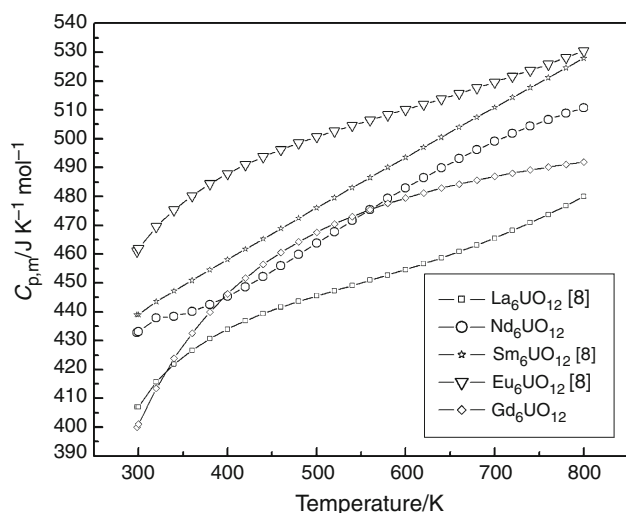


Fig. 8 Comparison of measured heat capacity data of $\text{RE}_6\text{UO}_{12}$ vs. temperature

$\text{Sm}^{3+} > \text{Eu}^{3+} > \text{Gd}^{3+}$ as higher the size of the ions, phonon frequencies are expected to be lower and therefore higher will be the heat capacity. From the Fig. 7, it is evident that the contribution of thermal expansion to heat capacity of $\text{RE}_6\text{UO}_{12}$ vary as $\text{Sm}^{3+} > \text{Nd}^{3+} > \text{Eu}^{3+} > \text{La}^{3+} \approx \text{Gd}^{3+}$. Though variation in the Schottky contribution to heat capacity for $\text{RE}_6\text{UO}_{12}$ is difficult to predict, it is expected that it should be least in the case of La^{3+} (No 4f electrons) and lesser in the case of Gd^{3+} (half filled 4f shell) than Nd^{3+} , Eu^{3+} and Sm^{3+} . The observed trend in the variation of heat capacity of $\text{RE}_6\text{UO}_{12}$ with varying RE could be due to the combined effect of all the three contributions to heat capacity mentioned above. The trend in the variation of heat capacity of RE_2O_3 as per the literature [23] is $\text{Gd} < \text{La} < \text{Nd} < \text{Sm} < \text{Eu}$. Whereas as per the present data the trend in the variation of heat capacity of $\text{RE}_6\text{UO}_{12}$ is same as that of RE_2O_3 except for that of La and Gd. Sahu et al. [24] measured the heat capacity of $\text{Gd}_6\text{UO}_{12}$ and compared with heat capacity data of $\text{RE}_6\text{UO}_{12}$ (RE = La, Sm and Eu) from our previous work [18] and observed that trend in the variation of heat capacity of $\text{RE}_6\text{UO}_{12}$ is same as the trend observed in that of RE_2O_3 . It has to be noted that in the case of RE_2O_3 , RE = La and Nd exhibits hexagonal structure [26] where as that of RE = Sm, Eu and Gd exhibits in both monoclinic and cubic form [26]. In the case of $\text{RE}_6\text{UO}_{12}$, all the compounds with RE = La to Gd exhibits rhombohedral structure. Therefore, trend in the variation in the heat capacity of $\text{RE}_6\text{UO}_{12}$ need not follow that of RE_2O_3 .

In our previous study [18] on the thermal expansion of $\text{RE}_6\text{UO}_{12}$, we could see an increasing trend in the % thermal expansion from the La through Nd to Sm compounds at any temperature. In the present study, we do not see the trend continuing to the Eu and Gd compounds. Thermal expansion of $\text{Eu}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ falls within

the range of that of the La, Nd and Sm systems as shown in Figs. 5, 6 and 7. This could be a manifestation of the peculiarity of the electronic configuration of the Eu and Gd ions. Gd^{3+} is $4f^7$ (half filled 4f level) and a part of Eu could be in $4f^7$, at high temperatures as result of the stability of Eu^{2+} state formed on loss of oxygen in $\text{Eu}_6\text{UO}_{12}$ [1].

From Figs. 5 and 6 it could be observed that there is a slight increase in the slope of the % expansion along the *a* axis vs. *T* plot and a slight decrease in slope of the % expansion along the *c* axis vs. *T* plot for $\text{Eu}_6\text{UO}_{12}$ at temperatures above 650 K. This is probably due to loss of oxygen from EuO_6 octahedra of $\text{Eu}_6\text{UO}_{12}$ at high temperatures and consequent stabilization of Eu^{2+} . Similar observation was reported by Berndt et al. [1].

Conclusions

Heat capacity of $\text{Nd}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ and thermal expansion coefficient of $\text{Eu}_6\text{UO}_{12}$ and $\text{Gd}_6\text{UO}_{12}$ were measured by DSC and HTXRD, respectively. The heat capacity of $\text{RE}_6\text{UO}_{12}$ varies in the order $\text{Eu}_6\text{UO}_{12} > \text{Sm}_6\text{UO}_{12} > \text{Nd}_6\text{UO}_{12} > \text{Gd}_6\text{UO}_{12} > \text{La}_6\text{UO}_{12}$. The trend is explained on the basis of the contributions of different factors namely phonons arising due to harmonic vibration, thermal expansion arising due to unharmonic vibration and Schottky effect due to excitation of electrons to higher energy levels. Though increasing trend in the % thermal expansion was observed for $\text{RE}_6\text{UO}_{12}$ for RE = La through Nd to Sm, the same is not maintained for Eu and Gd. Thermal expansion of $\text{RE}_6\text{UO}_{12}$ (RE = Eu and Gd) falls within the range of that of the La, Nd and Sm systems. This may be due to the stable electronic configuration $4f^7$ for Gd^{3+} and Eu^{2+} (formed partly due to loss of oxygen in $\text{Eu}_6\text{UO}_{12}$).

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